

# Lamellar Morphology of Poly(trimethylene terephthalate)/Poly(ether imide) Blends Studied via Small-Angle X-ray Scattering

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**ABSTRACT:** The lamellar morphology of a melt-miscible blend consisting of poly(trimethylene terephthalate) (PTT) and poly(ether imide) (PEI) prepared by solution precipitation has been investigated by means of optical polarized microscopy (POM) and small angle X-ray scattering (SAXS). From the observation under POM, it was suggested that PEI was predominantly segregated into the interlamellar and/or interfibrillar regions upon PTT crystallization since the PTT spherulitic morphologies of blends were volume-filling. From results of SAXS data analysis, a larger amorphous layer thickness was identified in the blends, showing that some PEI was incorporated inside the interlamellar regions

after crystallization. Despite the swelling of the amorphous layer, the amorphous layer thickness was relatively independent of the blend composition. It was concluded that amorphous PEI was located in the interlamellar regions of PTT as the weight fraction of PEI ( $w_{\text{PEI}}$ ) [ $\leq$ ] 0.1, while amorphous PEI was predominantly segregated into the interfibrillar regions of PTT as  $w_{\text{PEI}} > 0.1$ , and the extent of interfibrillar segregation increased with increasing  $w_{\text{PEI}}$ . © 2005 Wiley Periodicals, Inc. *J Appl Polym Sci* 99: 2421–2427, 2006

**Key words:** segregation morphology; poly(trimethylene terephthalate); poly(ether imide); interfibrillar segregation

## INTRODUCTION

Poly(trimethylene terephthalate) (PTT) is a linear aromatic polyester with three methylene moieties in the repeating unit. Although chemical structures are similar to poly(ethylene terephthalate) (PET) and poly(butylenes terephthalate) (PBT), the physical properties of PTT are much different from those of PET and PBT. PTT with excellent and economical properties for the applications such as films and fiber has been developed by Shell Chemical Company.<sup>1–5</sup> However, a lower glass transition temperature ( $T_g = 44^\circ\text{C}$ ) of PTT is a disadvantage in some applications. Since PTT is a crystalline polymer, its physical properties are expected to depend strongly on its crystalline morphology, which can be controlled by the crystallization kinetics. Blending the PTT with other polymers is an effective way for tuning the crystalline morphology. Poly(ether imide) (PEI) is a high-performance engineering plastic with a very high  $T_g$  ( $\sim 215^\circ\text{C}$ ).<sup>6</sup> PEI also possesses several advantageous properties such as ex-

cellent toughness and good electrical insulation, while PEI has a poor chemical resistance due to its amorphous nature. Therefore, it is expected that blending of PTT and PEI will offer an interesting route to combine the complementary properties of both polymers. Huang et al.<sup>6</sup> found that the PTT/PEI blends were miscible in the melt for the entire composition since the blends showed a single and composition-dependent  $T_g$  over the entire composition range and the crystallization kinetics also decreased with the addition of PEI. In the present study, we are interested in the investigation on the microstructure morphology of the PTT/PEI blends since the bulk physical properties are also stemmed from the microstructure morphology.

For a crystalline/amorphous system, the crystallization usually involves two types of polymer transport: diffusion of the crystalline component toward the crystal growth front and simultaneous segregation of the amorphous diluent away from the growth front. The diluent segregation can take place on three length scales, namely, interlamellar, interfibrillar, or interspherulitic. Interlamellar segregation occurs when the diluent is incorporated into the amorphous regions between individual lamellae. Interfibrillar segregation is characterized by the segregation of the amorphous diluent to the regions between the lamellar stacks,

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which consequently generates amorphous domains in between the lamellar stacks. Interspherulitic segregation is associated with the longest segregation distance, where the diluent is expelled out of the spherulites. Different morphological patterns may coexist in a given blend, leading to multiple locations for the amorphous diluent. These morphological patterns represent the diluent dispersion from the nanoscopic scale to the micrometer scale. Different scales of dispersion may lead to different properties.

In this article, we will investigate the microstructure morphologies of the PTT/PEI blends by means of small-angle X-ray scattering (SAXS). The effects of composition and thermal history on the segregation morphology of crystalline/amorphous state in the PTT/PEI blends will be discussed via SAXS data.

## EXPERIMENTAL

### Materials and sample preparation

PTT with  $M_n = 2.1 \times 10^4$  and  $M_w = 4.63 \times 10^4$  was obtained from Shinkong Synthetic Fibers Co. Taiwan. PEI with  $M_w = 1.2 \times 10^4$  and  $M_w = 3.0 \times 10^4$  was acquired from General Electrical (GE, Ultem 1000) Co. America. PTT were blended with PEI by solution precipitation. A series of PTT/PEI blends of designated weight composition ratios were stirred and dissolved in dichloroacetic acid at 80°C. These blends were subsequently recovered by precipitation in 10-fold excess volume of water. The blends were thoroughly washed with a large amount of water and then dried in a vacuum oven at 105°C for 3 days.

Samples for SAXS study were prepared by pressing the blends between two pieces of Teflon films on a Linkam HFS91 hot stage at  $270 \pm 0.2^\circ\text{C}$  for 3 min. The samples were subsequently transferred directly into an oven equilibrated at the desired crystallization temperature ( $T_c$ ) to allow the crystallization of PTT for 5 h. Optical microscopy confirmed that volume-filling spherulitic samples were obtained through such a crystallization condition.

### Polarized optical microscopy

The spherulitic morphology and growth were monitored with a Zeiss polarized optical microscope. The sample was first melted on a Linkam HFS91 hot stage at  $270 \pm 0.2^\circ\text{C}$  for 3 min, and then quickly transferred to another hot stage equilibrated at the desired crystallization temperature ( $T_c$ ), where the spherulitic growth was monitored. Micrographs were taken at intervals for measuring the spherulite radii ( $R$ ) at various time periods. The growth rate was calculated from the change of spherulitic radius with time,  $dR/dt$ .

### SAXS measurement

The SAXS experiments were conducted at the Department of Chemical Engineering, National Tsing Hua University, Taiwan. SAXS measurements were performed using a Bruker NanoSTAR SAXS instrument. The X-ray source, a 1.5 kW X-ray generator (Kristalloflex 760) equipped with a Cu tube, was operated at 35 mA and 40 kV. The scattering intensity was detected by a two-dimensional position-sensitive detector (Bruker AXS) with  $512 \times 512$  channels. All data were corrected by the empty beam scattering, the sensitivity of each pixel of the area detector. The area scattering pattern has been radially averaged to increase the photon counting efficiency compared with the one-dimensional linear detector. The intensity profile was output as the plot of the scattering intensity ( $I$ ) versus the scattering vector,  $q = 4\pi/\lambda \sin(\theta/2)$ , where  $\theta$  is the scattering angle and  $\lambda$  is the X-ray wavelength.

### Bulk crystallinity measurements

Bulk crystallinities of crystallized PTT/PEI blends were calculated from the enthalpy of melting ( $\Delta h_f$ ). The sample used for  $\Delta h_f$  measurement was cut directly from the specimens for SAXS measurement. The enthalpy of melting was measured by a METTLER TOLEDO DSC822e differential scanning calorimeter (DSC). The bulk crystallinity ( $\phi_c$ ) was calculated by taking 146 J/g as the enthalpy of melting of 100% crystalline PTT,<sup>3,7</sup> i.e.,  $\phi_c = \Delta h_f / 146$ .

## RESULTS AND DISCUSSION

### Spherulitic morphology

The spherulitic morphology of crystalline polymers can be readily observed by POM. Figure 1 shows the spherulitic morphologies of PTT/PEI blends crystallized at  $T_c = 190^\circ\text{C}$ . PTT spherulitic growths were observed to be linear with time and no apparent evidence of liquid–liquid phase separation was found up to the point of spherulite impingement for blend compositions studied. The observed spherulitic morphology was volume-filling, implying that the PTT/PEI blends mainly created the intraspherulitic segregation (i.e., interlamellar or interfibrillar segregation).

### Lamellar morphology

The lamellar microstructure of a crystalline/amorphous polymer blend can conveniently be probed from the profile of scattering intensity ( $I(q)$ ) with scattering vector ( $q$ ) through the SAXS measurement. In the present study, we utilized the SAXS technique to probe the lamellar morphologies of PTT/PEI blends. Figure 2 depicts the Lorentz-corrected SAXS profiles of PTT/PEI blends crystallized at  $T_c = 200, 180, 160^\circ\text{C}$ ,

respectively. In Figure 2, we observed the scattering peak ( $q_{\max}$ ) associated with the electron density contrast between the alternating crystalline and amorphous layers at each blend composition. The  $q_{\max}$  shifted to a smaller position with the addition of 10 wt % PEI content and then maintained seemingly unchanged with further increasing PEI composition, indicating the value of long period ( $L^{\text{Bragg}}$ ) calculated by Bragg's law ( $L^{\text{Bragg}} = 2\pi/q_{\max}$ ) for blends at each  $T_c$  is remained constant but its value is larger than that of neat PTT.

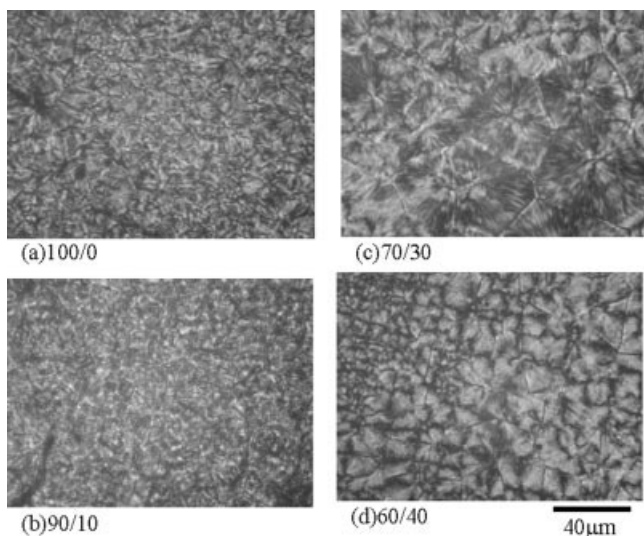
The morphological parameters including the crystal thickness ( $l_c$ ), amorphous layer thickness ( $l_a$ ), and the long period ( $L = l_c + l_a$ ) may be evaluated by the one-dimensional correlation function ( $\gamma(z)$ )<sup>8</sup> defined to:

$$\gamma(z) = \frac{1}{Q} \int_0^{\infty} q^2 I(q) \cos(qz) dq \quad (1)$$

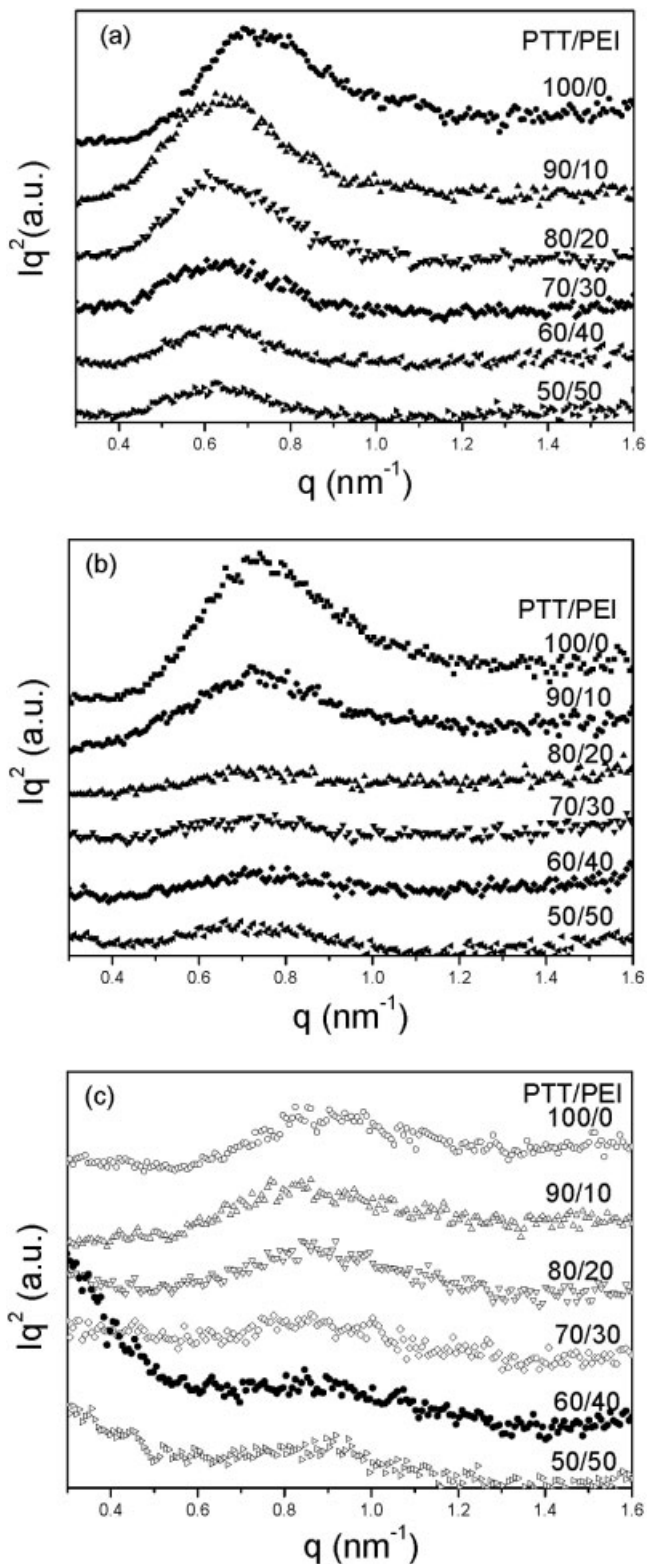
where  $z$  is the direction along which the electron density is measured.  $Q$  is the scattering invariant.

$$Q = \int_0^{\infty} I(q) q^2 dq \quad (2)$$

Since the experimentally accessible  $q$  range is finite, extrapolation of intensity to both low and high  $q$  is necessary for the integrations. Extrapolation to zero was accomplished by the Debye-Bueche model<sup>9</sup>



**Figure 1** Spherulitic morphologies of PTT/PEI blends crystallized at  $T_c = 190^\circ\text{C}$  under POM.



**Figure 2** Lorentz-corrected SAXS profiles of PTT/PEI blends crystallized at (a)  $T_c = 200^\circ\text{C}$ ; (b)  $T_c = 180^\circ\text{C}$ ; (c)  $T_c = 160^\circ\text{C}$ .

$$I(q) = \frac{A}{(1 + a_c^2 q^2)^2} \quad (3)$$

where  $A$  is a constant and  $a_c$  is the correlation length.  $A$  and  $a_c$  can be determined from the plot of  $I(q)^{-0.5}$  versus  $q^2$  using the intensity data at low  $q$  region. Extension to large  $q$  can be performed using the Porod–Ruland model<sup>10–11</sup>

$$I(q) = K_p \frac{\exp(-\sigma^2 q^2)}{q^4} + I_{fl} \quad (4)$$

where  $K_p$  is the Porod constant,  $\sigma$  is a parameter related to the thickness of interphase between crystal and amorphous layers, and  $I_{fl}$  is the background intensity arising from thermal density fluctuation. The values  $K_p$ ,  $\sigma$ , and  $I_{fl}$  were obtained by curve fitting the intensity profile at high  $q$  region. Figure 3 shows the one-dimensional correlation functions of PTT/PEI blends crystallized at various  $T_c$ s. Assuming the corresponding two-phase model,  $l_c, l_a$ , the most probable value of  $L$  can be estimated via simple geometric analysis of  $\gamma(z)$ , as demonstrated in Figure 4. The thickness of the thinner layers ( $l_1$ ) is given by the intersection between the straight line extended from the self-correlation triangle and the baseline given by  $-A^8$ . The average thickness of the thicker layer is then obtained from  $l_2 = L - l_1$ . The assignment of  $l_1$  and  $l_2$  is governed by the magnitude of the linear crystallinity ( $\phi_c^{lin}$ ). When  $\phi_c^{lin} < 0.5$ , the crystals contribute the smaller thickness; then  $l_1 = l_c$  and  $l_2 = l_a$ . The inverse is true for  $\phi_c^{lin} > 0.5$ . Linear crystallinity,  $\phi_c^{lin}$  is defined as

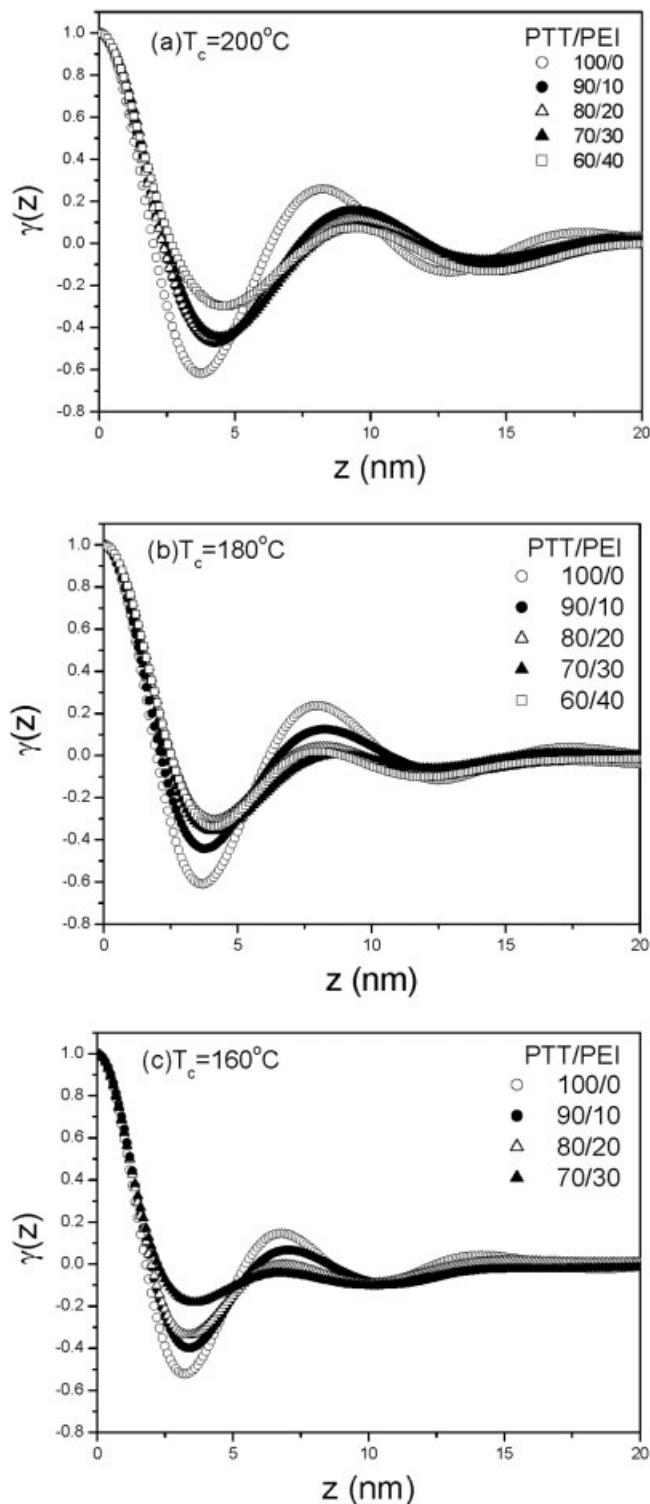
$$\phi_c^{lin} = \frac{l_c}{L} = \frac{l_c}{l_c + l_a} \quad (5)$$

Provided that spherulites are volume-filling,  $\phi_c^{lin}$  is related the bulk crystallinity,  $\phi_c$  by

$$\phi_c = \phi_s \phi_c^{lin} \quad (6)$$

where  $\phi_s = \phi_c / \phi_c^{lin}$ , the volume fraction of lamellar stacks in the sample. Since  $\phi_s < 1$ , eq. (6) prescribes that the bulk crystallinity cannot be higher than the linear crystallinity. As a result, the assignment of  $l_1$  and  $l_2$  would be rather straightforward if the bulk crystallinity is greater than 0.5, because in this case the larger length,  $l_2$ , must correspond to the crystal thickness and  $l_1$  to the amorphous layer thickness. In the present study, for neat PTT and all blend compositions at each  $T_c$ ,  $\phi_c$  is below 0.5 as shown in Figure 6, it is reasonable to assign  $l_1 = l_c$  and  $l_2 = l_a$ .

Figure 5 shows the variations of  $L$ ,  $l_c$  and  $l_a$  with the weight fraction of PEI ( $w_{PEI}$ ) at various  $T_c$ s. It is observed in Figure 5(a) that the  $L$  value increased as the addition of 10 wt % PEI content but did not vary with further increasing PEI content. This result is consistent with the observation in Figure 2. From the observation



**Figure 3** One-dimensional correlation functions of PTT/PEI blends crystallized at (a)  $T_c = 200^\circ\text{C}$ ; (b)  $T_c = 180^\circ\text{C}$ ; (c)  $T_c = 160^\circ\text{C}$ .

in Figure 5(b), it is obvious that the  $l_c$  value increased with elevating  $T_c$ . According to the secondary nucleation theory,<sup>12–13</sup> the initial crystal thickness is given by

$$l_g^* = \frac{2\sigma_e T_m^0}{\Delta h_f^0 (T_m^0 - T_c)} + \delta l \quad (7)$$

where  $T_m^0$  is the equilibrium melting point,  $\sigma_e$  is the fold surface free energy,  $\Delta h_f^0$  is the bulk enthalpy of melting per volume, and  $\Delta T$  is defined as the degree of undercooling (i.e.,  $\Delta T = T_m^0 - T_c$ ). At low to moderate degree of undercooling,  $\delta l$  is small, eq. (7) thus reduces to

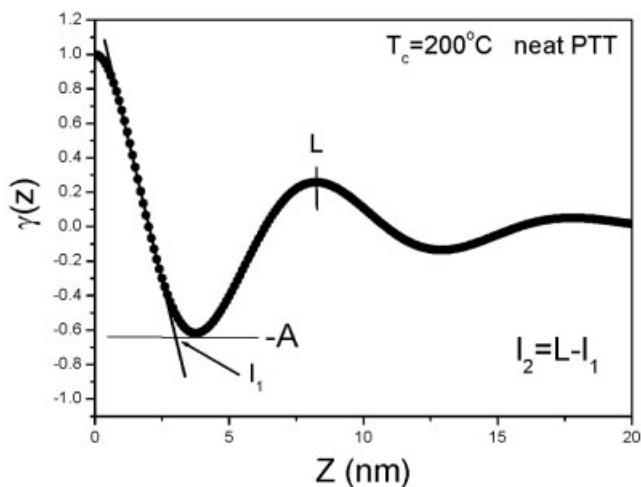
$$l_g^* = \frac{2\sigma T_m^0}{\Delta h_f^0 (T_m^0 - T_c)} \quad (8)$$

The final crystal thickness, according to the notation of Hoffman and Weeks,<sup>12–13</sup> is  $\gamma$  times the initial thickness, i.e.,

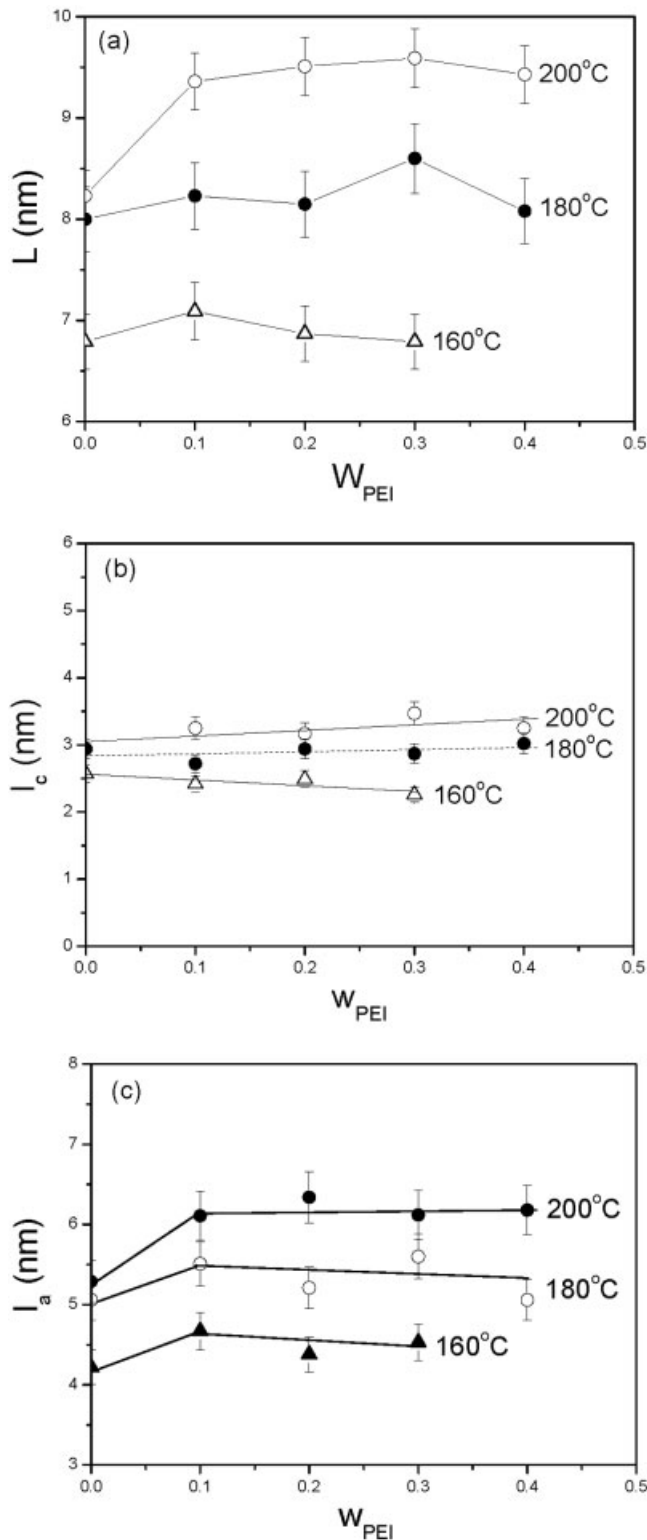
$$l_c = \gamma l_g^* \quad (9)$$

where  $\gamma$  is the so-called “lamellar thickening factor”. Equation (8) prescribes that the  $l_g^*$  is inversely proportional to the degree of undercooling. Therefore, the higher the  $T_c$ , the lower the degree of undercooling (i.e.,  $\Delta T = T_m^0 - T_c$ ), leading to thicken the crystal thickness. Such a consequence, the  $l_c$  value increased with elevating  $T_c$  in Figure 5(b), is attributed to the depression of undercooling.

In addition, it is worth noting in Figure 5(b) that the  $l_c$  appreciably increased for both  $T_c = 200$  and  $180^\circ\text{C}$  but decreased for  $T_c = 160^\circ\text{C}$  with the increase of PEI composition. In other words, upon blending, the crystal thickness of PTT for both  $T_c = 200$  and  $180^\circ\text{C}$  increased but that for  $T_c = 160^\circ\text{C}$  reduced. As described above, the decrease of  $T_m^0$  would lead to the larger  $l_g^*$  so as to form thicker crystals. The increase in



**Figure 4**  $l_c$ ,  $l_a$ , the most probable value of  $L$  can be estimated via simple geometric analysis of  $\gamma(z)$  (the thinner length,  $l_1$ , has been assigned to be  $l_c$  and the larger length,  $l_2$ , to be  $l_a$ ).



**Figure 5** Variations of  $L$ ,  $l_c$ , and  $l_a$  with PEI composition (a)  $L$ ; (b)  $l_c$ ; (c)  $l_a$ .

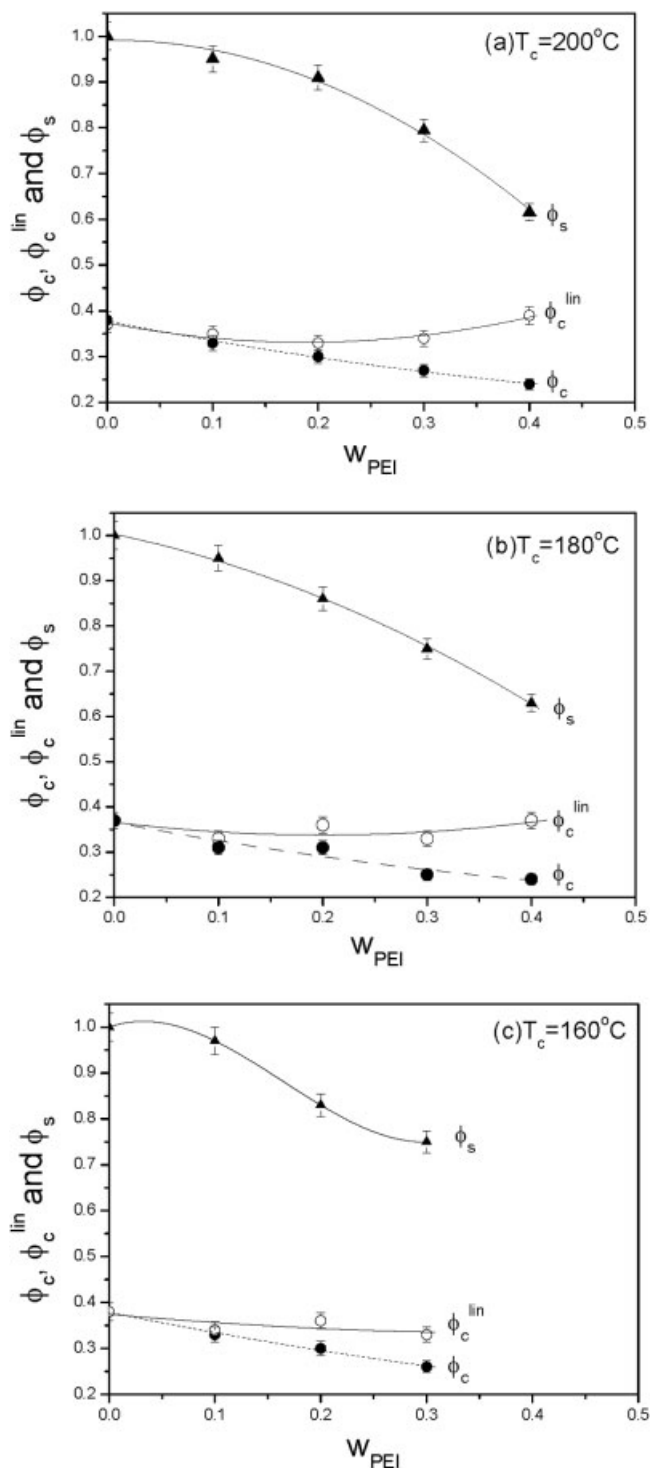
$l_c$  with increasing amorphous component has been observed as reported in literature,<sup>14–16</sup> attributed to the depression in  $T_m^0$  upon blending at a given  $T_c$ . Therefore, in Figure 5(b) the  $l_c$  value increased with

increasing PEI composition for  $T_c = 200$  and  $180^\circ\text{C}$ , which may be ascribed to the decrease in  $T_m^0$ . Above discussion concerning the increase in  $l_c$  has mainly focused on the consideration of crystallization thermodynamics. Nevertheless, the sizes of crystals may be also affected by the crystallization kinetic factor, which may lead to the maintaining of a constant thickness of crystals even with increasing amorphous component such as *i*-polystyrene (*i*PS)/*a*-polystyrene (*a*PS),<sup>17</sup> sPS/*a*PS<sup>18</sup> and *i*PS/poly(propylene oxide) (PPO)<sup>19</sup> blends and so on. In Figure 5(b), the slight decrement of  $l_c$  with increasing PEI for  $T_c = 160^\circ\text{C}$  could be also considered due to both the crystallization kinetics or morphological factors resulted.

Figure 5(c) displays the effect of PEI composition on the thickness of the amorphous layers. The trend of  $l_a$  with PEI composition for three  $T_c$ s is quite similar:  $l_a$  increased as  $w_{\text{PEI}} [\leq] 0.1$  but did not vary with further increasing PEI composition. Despite the swelling of amorphous layers, the thickness of amorphous layers was relatively independent of the PEI composition, where all blends exhibited approximately the same  $l_a$ . Because a complete interlamellar incorporation should result in a monotonic increase of amorphous layer thickness with increasing PEI composition, the observation from Figure 5(c) implies an incomplete incorporation of PEI into the interlamellar regions. In other words, some PEI molecules were expelled out of the interlamellar regions. A similar phenomenon has also been observed in PET/PEI blends.<sup>20</sup>

The extent of interfibrillar segregation can also be determined from the magnitude of the volume fraction of lamellar stacks (e.g., the value of  $\phi_s$ ). The magnitude of  $\phi_s$  is closely connected with the morphological structure. In the case of complete interlamellar segregation, the whole sample is homogeneously filled with lamellar stacks, so  $\phi_s = 1$ . A smaller  $\phi_s$  implies a larger extent of interfibrillar segregation. Figure 6 depicts the variations of  $\phi_c$ ,  $\phi_c^{\text{lin}}$ , and  $\phi_s$  with PEI composition.  $\phi_s$  remained essentially constant at approximately unity when  $w_{\text{PEI}} [\leq] 0.1$ . In this case, PEI molecules were predominantly incorporated into the interlamellar regions. When  $w_{\text{PEI}} > 0.1$ , the  $\phi_s$  obviously dropped with PEI composition, indicating the existence of extralamellar segregation after PTT crystallization, and the extent of extralamellar segregation increased with increasing of PEI composition.

In addition, in Figure 2(c) a strong intensity upon at the low- $q$  region can be observed for blends of 60/40–50/50 blends crystallized at  $T_c = 160^\circ\text{C}$ . The abrupt intensity at the low- $q$  region is termed “zero-angle scattering.” The observed scattering patterns on the 60/40–50/50 blends crystallized at  $T_c = 160^\circ\text{C}$  are characterized by the superposition of a monotonically decayed profile and a scattering peak associated with PTT lamellar stacks. The large low- $q$  intensity or



**Figure 6** Variations of  $\phi_c$ ,  $\phi_c^{\text{lin}}$  and  $\phi_s$  with PEI composition (a)  $T_c = 200^\circ\text{C}$ ; (b)  $T_c = 180^\circ\text{C}$ ; (c)  $T_c = 160^\circ\text{C}$ .

“zero-angle scattering” signifies the presence of heterogeneity having a size larger than the crystalline and amorphous layers. According to Schultz,<sup>21</sup> the zero-angle scattering may be attributed to a large (greater than lamellar length scales) individual amorphous domain inserted into the stacks of several lamellae to

create a gap in the lamellar stacks. The electron density contrast between this large amorphous domain and the stacking of alternating crystalline–amorphous layers gives rise to the zero-angle scattering. In PTT/PEI blends, the larger amorphous domains exterior to the PTT lamellar stack referred to the PEI molecules, which are located in the interfibrillar regions. These exterior amorphous domains were generated as PEI was segregated beyond several layers of lamellae. This type of zero-angle scattering has also been observed in PET/PEI blends,<sup>20</sup> where the morphology was induced by a combined crystallization and liquid–liquid demixing.

It should be noted that zero-angle scattering does not necessarily appear in the Lorentz-corrected profiles even if interfibrillar segregation is evident. If the interfibrillar segregation distance is so large (e.g., approaches the order of micrometers) that the amorphous diluent is expelled beyond many layers of lamellae, the interference between the resultant large lamellar stack domains and interfibrillar amorphous domains could generate the zero-angle scattering at the  $q$  region that is too small to be accessed by the typical SAXS instrument. In the present study, for the blends crystallized at  $T_c = 200$  and  $180^\circ\text{C}$ , the amorphous PEI was expelled beyond many layers of PTT lamellae, the interference between the resultant large lamellar stack domains and interfibrillar amorphous domains of amorphous PEI could generate the zero-angle scattering at the  $q$  region that is too small to be accessed by the typical SAXS instrument. Therefore, it may be suggested that the distance of PEI segregation at  $T_c = 160^\circ\text{C}$  is shorter than those at  $T_c = 200$  and  $180^\circ\text{C}$ .

## CONCLUSIONS

The lamellar morphology of PTT/PEI blends has been investigated by means of POM and SAXS. From the observation under POM, it was suggested that PEI was predominantly segregated into the interlamellar and/or interfibrillar regions upon PTT crystallization since the PTT spherulitic morphologies of blends were volume-filling. On the basis of the composition variation of  $l_a$ , the magnitude of  $\phi_s$  revealed that amorphous PEI was located in the interlamellar regions of PTT as  $w_{\text{PEI}} [\leq] 0.1$ , while amorphous PEI was predominantly segregated into the interfibrillar regions of PTT as  $w_{\text{PEI}} > 0.1$ , and the extent of interfibrillar

segregation increased with increasing  $w_{\text{PEI}}$ . The zero-angle scattering can be observed only for 60/40 and 50/50 blends crystallized at  $T_c = 160^\circ\text{C}$  but not for those crystallized at  $T_c = 200$  and  $180^\circ\text{C}$  from the Lorentz-corrected SAXS profiles. This may be attributed to the fact that amorphous PEI at  $T_c = 200$  and  $180^\circ\text{C}$  was segregated beyond many layers of PTT lamellae, the interference between the resultant PTT large lamellar stack domains and interfibrillar amorphous domains of PEI could generate the zero-angle scattering at the  $q$  region that is too small to be accessed by the typical SAXS instrument. Therefore, it may be suggested that the distance of PEI segregation at  $T_c = 160^\circ\text{C}$  is shorter than those at  $T_c = 200$  and  $180^\circ\text{C}$ .

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